

Kochi 07/12/2011

Discretizing the One-dimensional KPZ Equation

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- KPZ equation / exact solution

- lattice discretization of KPZ

jointly with T. Sasamoto (2009)

- lattice discretization of SHE

jointly with S. Prokhorov (2011)

1. KPZ equation / exact solution

Kardar, Parisi, Zhang (1986)

height function $h(t): \mathbb{R} \rightarrow \mathbb{R}, t \geq 0$

$$\frac{\partial}{\partial t} h = \frac{1}{2} \left(\frac{\partial}{\partial x} h \right)^2 + \frac{1}{2} \frac{\partial^2}{\partial x^2} h + W$$

- $W(x,t)$ Gaussian white noise $\langle \cdot \rangle$
- Cauchy problem, initial data $h(x,0)$ could be random

⇒ Cole-Hopf transform

$$Z(x,t) = e^{h(x,t)}$$

$$\frac{\partial}{\partial t} Z = \frac{1}{2} \frac{\partial^2}{\partial x^2} Z + \underbrace{WZ}$$

Ito \uparrow // stochastic heat equation //

auxiliary Brownian motion $B(t)$ "directed polymer"

$$Z(x, t) = \mathbb{E}_x \left(\exp \left[\int_0^t W(B(s), t-s) ds \right] Z(B(t), 0) \right)$$

// random partition function //

$h = \log Z$ // random free energy //

$$Z > 0, \quad \langle Z(x, t)^2 \rangle < \infty$$

Bertini, Giacomin (1995)

$Z(x, t)$ Hölder $\frac{1}{2}$ in x , Hölder $\frac{1}{2}$ in t

→ Cole-Hopf solution

$$h = \log Z$$

Hairer (2011) solution theory for KPZ, method of rough paths

- can be used for

1) $\frac{1}{2} \left(\frac{\partial h}{\partial x} \right)^2 + \frac{1}{4} \left(\frac{\partial h}{\partial x} \right)^4$ renormalizes to $\lambda \left(\frac{\partial h}{\partial x} \right)^2$

2) vector-valued KPZ

1D nonlinear fluctuating hydrodynamics
van Beijeren 2011

$$\frac{\partial}{\partial t} h_\alpha = c_\alpha \frac{\partial}{\partial x} h_\alpha + \frac{1}{2} \sum_{\beta, \gamma} g_\alpha^{\beta\gamma} \frac{\partial}{\partial x} h_\beta \frac{\partial}{\partial x} h_\gamma + \underbrace{\frac{1}{2} \frac{\partial^2}{\partial x^2} h_\alpha + W_\alpha}_{\text{"more general"}}$$

Cole-Hopf solution is **NOT** available

exact solution

Sharp wedge initial data

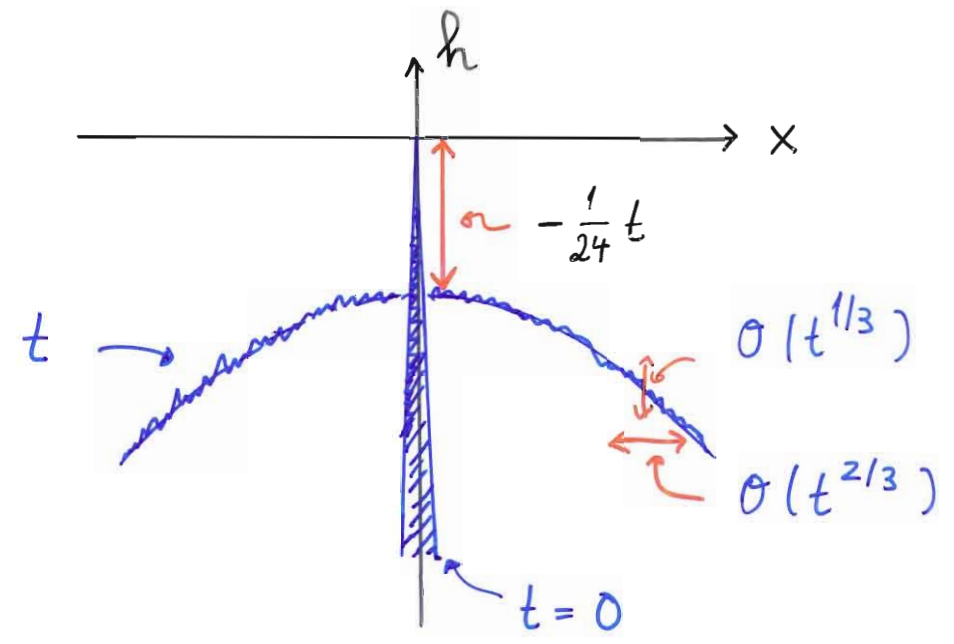
$$z(x, 0) = \delta(x)$$

$$h(x, t) = -\frac{1}{24}t + \left(\frac{1}{2}t\right)^{1/3} \xi(t)$$

$\mathbb{P}(\xi(t) \in du) = \text{Fredholm determinant}$
on $L^2(\mathbb{R})$

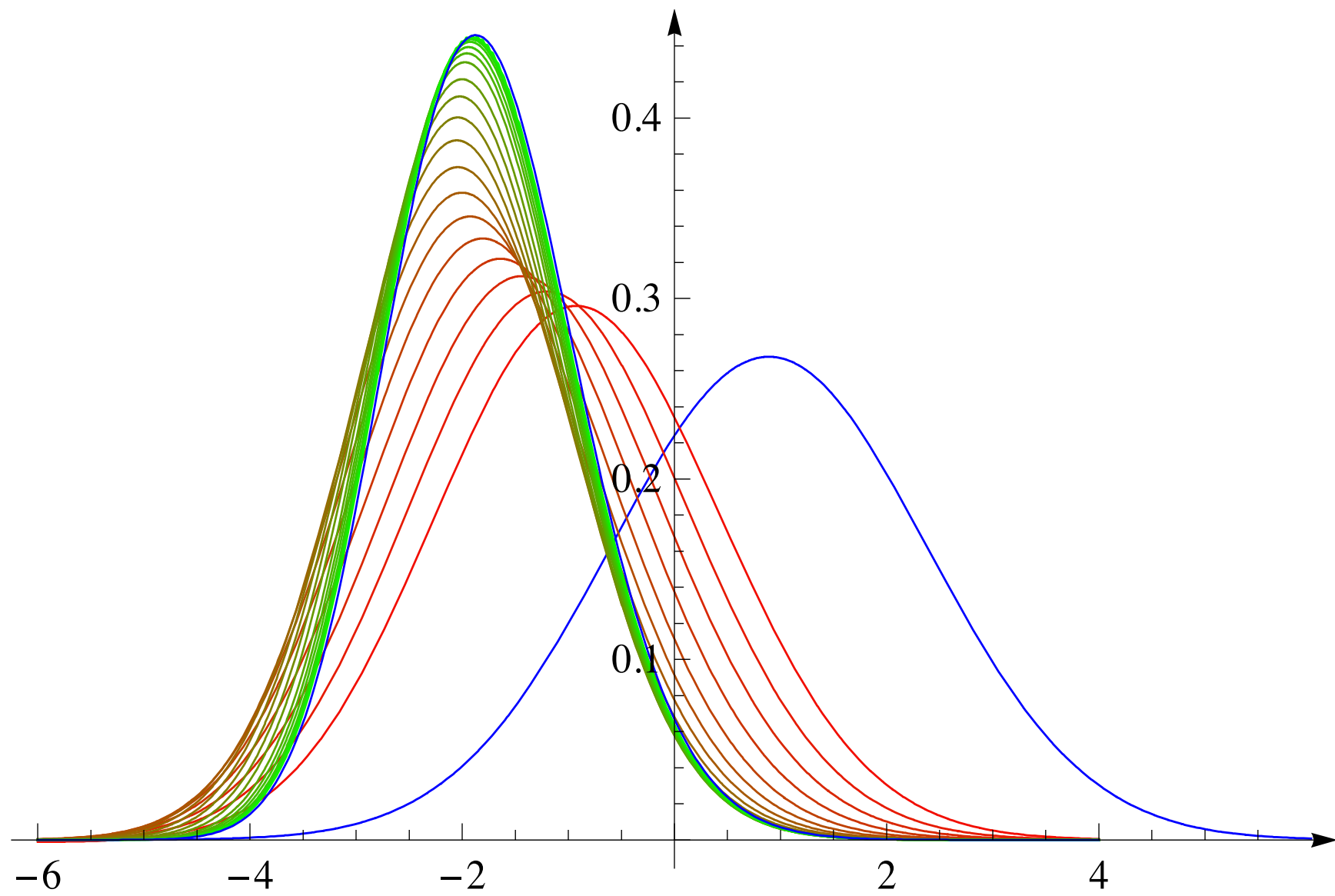
plot: S. Prolhac (2011)

based on the method of Bornemann (2008)



fixed t : $x \mapsto h(x, t) + \frac{1}{2t}x^2$

stationary in x



2. Lattice approximation for KPZ

1D conservation law: $u = \frac{\partial}{\partial x} h$, $\frac{\partial}{\partial t} u = \frac{\partial}{\partial x} u^2 + \frac{1}{2} \frac{\partial^2}{\partial x^2} u + \frac{\partial}{\partial x} W$

- discretize $\mathbb{R} \rightarrow \mathbb{Z}$ coupled diffusions

$u(x, t) \rightarrow u_j(t)$ • $W \rightsquigarrow dB_j(t)$ independent Brownian motions

• $\frac{\partial^2}{\partial x^2} \rightsquigarrow$ difference Laplacian Δ

• $\frac{\partial}{\partial x} u^2$? local // invariant measure

$$\rightarrow du_j = \left(\lambda (u_{j+1}^2 + u_j u_{j+1} - u_{j-1} u_j - u_{j-1}^2) + (\Delta u)_j \right) dt + dB_j - dB_{j-1}$$

invariant measures: i.i.d. Gaussians

one-parameter family: mean

link to PASEP

generator L

invariant measure, expand in Hermite polynomials $\{n_j, j \in \mathbb{Z}\}$, $n_j \in \mathbb{N}$

P projects on subspace $\{n_j = 0, 1, j \in \mathbb{Z}\}$

→ PLP is generator of PASEP ←

variance bounds based on Landim, Quastel, Salmhofer, Yau (2004)

- stationary process, mean 0

$$t^{5/4} \leq \sum_{j \in \mathbb{Z}} \mathbb{E} (u_0(0) u_j(t))^2 \leq t^{3/2}$$

|| should be $t^{4/3}$ ||

3. Lattice approximation for SHE

$$Z(x, t) \rightsquigarrow Z_j(t) \quad dZ_j = (\Delta Z)_j dt + Z_j db_j, \quad j \in \mathbb{Z}$$

sharp wedge:

$$Z_j(0) = \delta_{j0}$$

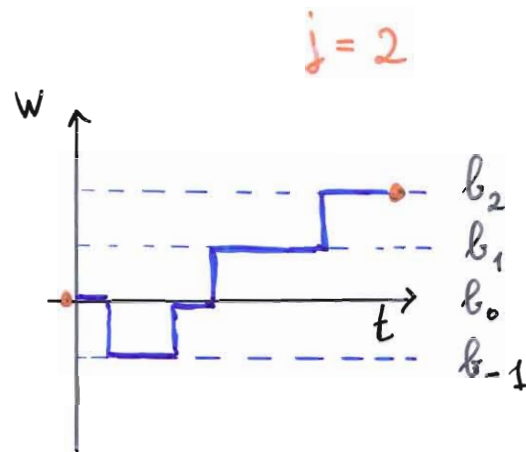
symmetric random walk $w(t)$:

$$Z_j(t) = \mathbb{E}_0 \left(e^{\int_0^t dw(s) \delta_{w(s), j}} \delta_{w(t), j} \right)$$

$$\rightsquigarrow dh_j = (e^{h_{j+1} - h_j} + e^{h_{j-1} - h_j} - 2) dt + db_j$$

$$u_j = h_j - h_{j-1}$$

$$du_j = (e^{u_{j+1}} + e^{-u_j} - e^{u_j} - e^{-u_{j-1}}) dt + db_j - db_{j-1}$$



point-to-point
directed polymer

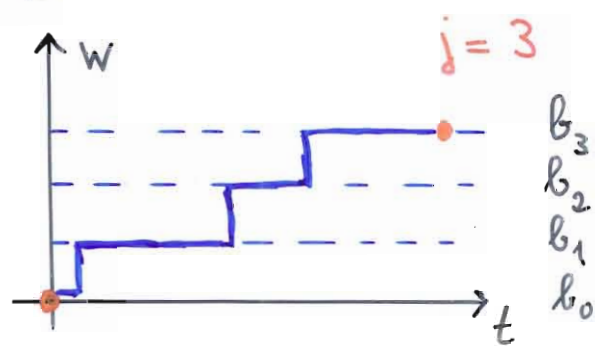
? invariant measure ?

⇒ asymmetric discretization

$$dZ_j = Z_{j-1} dt + Z_j db_j, \quad Z_j(0) = \delta_{j0}, \quad j \in \mathbb{Z}$$

asymmetric random walk $w(t)$

$$Z_j(t) = E_0 \left(e^{\int_0^t db_w(s)} \delta_{w(t); j} \right) e^t$$



|| zero temperature limit
|| maximize random energy

Baryshnikov (2001)
O'Connell, Yor (2001)

RMT

$$dh_j = e^{h_{j-1} - h_j} dt + db_j$$

$$u_j = h_{j+1} - h_j$$

$$du_j = (e^{-u_j} - e^{-u_{j-1}}) dt + db_{j+1} - db_j$$

⌋ stationary measures:

parameter $r > 0$

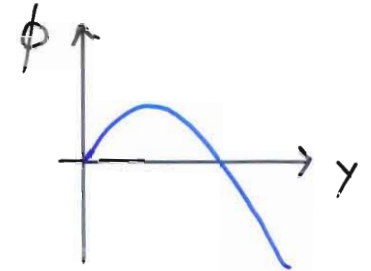
product $\frac{1}{\Gamma(r)} e^{-e^{-x} - rx} dx \quad \langle \cdot \rangle_r$

macroscopic shape

density $\rho = \langle u_0 \rangle_T = -\Psi(\tau)$

Digamma function $\Psi = \frac{\Gamma'}{\Gamma}$

current $j = -\langle e^{-u_0} \rangle_T = -\tau$



shape function $\phi(\gamma) = \inf_{\rho} (-j(\rho) - \rho\gamma), \gamma \geq 0$

law of large numbers

$\lim_{t \rightarrow \infty} \frac{1}{t} \ln h_{L\gamma t}(t) = \phi(\gamma) \quad (*)$

Theorem (O'Connell, Moriarty, 2007)

$\lim_{n \rightarrow \infty} \frac{1}{n} \log h_n(kn) = f(k)$

$f(k) = \inf_{s \geq 0} (ks - \Psi(s))$

in agreement with (*)

$\gamma k = 1$

fluctuations

KPZ scaling theory

Krug, Meakin, Halpin-Healy, 1992

$$\lim_{t \rightarrow \infty} \frac{1}{t^{1/3}} (h_{[yt]}(t) - t\phi(y)) = \left(\frac{1}{2} \lambda A^2\right)^{1/3} \overset{\text{TW}}{\sim} \quad (*)$$

GUE Tracy-Widom r.v.

- fix y , $\rho = \phi'(y)$

- $\lambda(\rho) = -j''(\rho)$

- $A = \langle u_0^2 \rangle_{\tau} - \langle u_0 \rangle_{\tau}^2 \quad A(\tau) = \psi'(\tau)$

Theorem (Borodin, Corwin, 2011)

• Theorem 5.2.12 of Macdonald processes

$$k > k^* > 0$$

$$\lim_{n \rightarrow \infty} \frac{1}{n^{1/3}} (h_n(kn) - n f(k)) = \left(\frac{1}{2} \frac{1}{(-f''(k))}\right)^{1/3} \overset{\text{TW}}{\sim}$$

in agreement with (*)

Conclusions

- KPZ limit of semi-discrete directed polymer
- work directly with KPZ equation

$$\beta \rightarrow 0, n, t \rightarrow \infty$$
$$\beta^4 n = 1$$
$$\beta^4 t = 1$$

ONLY replica



THANK YOU

for a really wonderful workshop!

